## **Magnonic Neural Networks**

Spin Wave Inputs

# Nonlinear Chiral Magnonic Resonators: Towards Magnonic Neurons

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## **Chiral Magnonic Resonators**

### **Magnonic Neurons:**

- Feed-forward information propagation -> Non-reciprocal transmission
- Activation -> Non-linear transmission

We propose **chiral magnonic resonators**  to act as the artificial magnonic neurons.





# **A Stripe Resonator and Thin Film Methodology**

 $\widehat{\Xi}^{18}$ 

 $\overline{\bigcirc}$  16

 $\sum 14$ 

Can an all magnonic neural network be built?

 $\frac{12}{\text{Hz}}$  12

Magnonic Spin Wave Outputs Neurons



 $B_{bias} = 100$  mT

(film only)

Chiral coupling arises from the relative sense of rotation between a resonator's precession and a spin wave's dynamic stray field:

*What are chiral magnonic resonators?*

What are the properties of a magnonic neuron?

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## **Abstract**

described in terms of a nonlinear shift of the resonant frequency ('detuning'), which results in amplitude-dependent transmission of monochromatic spin waves. This behaviour can be harnessed to realise a sigmoid-like In this work, we explore chiral magnonic resonators as building blocks of artificial neural networks. Using micromagnetic simulations and analytical modelling, we demonstrate that the first anti-symmetric confined ('dark') mode of a stripe chiral magnonic resonator may exhibit a strongly nonlinear response when resonantly excited by incoming spin waves, owing to energy concentration. For modest excitation levels, the effect can be activation, and thus implement artificial neurons in a network linked by spin waves propagating in a linear medium. The nonlinearity is manifested in bistability and hysteresis akin to those occurring in non-linear oscilla when the excitation strength exceeds a threshold set by the decay rate of the mode. In magnonic resonators, the latter includes both the Gilbert damping and the radiative decay due to the coupling with the medium. The results of our simulations are well described by a phenomenological model in which the nonlinear detuning of the confined mode is quadratic in its amplitude, while the propagation in the medium is linear.

In the linear case we see the curves We now see the transmission curves  $\|\cdot\|_{M_{\alpha}}^{\frac{A_0}{M_{\alpha}}=0.06}$ for the transmission for waves for increasing amplitudes of incident propagating towards the right and the spin waves. At low powers (blue and left for the two resonant modes. In black), the transmission experiences both cases nearly identical behaviour minimal frequency shifting. At for right-propagating spin waves moderate powers (light blue) there is occurs. A sharp Lorentzian dip with more detuning, and the curves are maximum attenuation of about 10%. now asymmetric. At higher powers Ultimately, the character of the (green to pink), the left edge of the resonant modes differs but the curves have become vertical, and the transmission behaviour is very curves detuned further. Beyond the similar. We may probe the nonlinear dashed line, the critical region of behaviour by looking only at dark bistability is entered. Indicated by the mode resonance. dashed curves, the transmission has developed bistable behaviour.



### **Transmission Activation Modelling**

Additionally, if we drive the resonator at a frequency below the resonance, we create a power limiter (top right). The shifting of the transmission curves has now aligned the minima of transmission at the lower frequency. The snapshot shows the reduction of spin wave amplitude in transmission.



We launch a broadband wave packet:

We can model the local resonator mode  $\phi$  as a driven, damped harmonic oscillator with cubic nonlinearity:

We compute the transmission coefficient as the ratio of the complex Fourier amplitudes, averaged in the  $yz$  crosssection:



*What is the frequency dependence of*  $T(\omega)$ ?



# 50 nm

### Wave number (rad/ $\mu$ m)

We study the transmission of a stripe chiral magnonic resonator spaced to a thin film waveguide. We assume permalloy parameters throughout and apply a bias field locally to the film to saturate the magnetisation in the backward volume geometry. The resonator is saturated by its shape anisotropy. The resonator possesses two resonant modes. A lower frequency quasi-uniform mode with approximately uniform amplitude and phase, and a higher frequency mode. This higher frequency mode is the first anti-symmetric or 'dark' mode of the resonator.

*How are spin waves transmitted across?*



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response (top left). The shifting of the transmission curve to lower frequencies has allowed spin waves at higher amplitudes to break through. The snapshot of the dynamic magnetisation shows this breakthrough of the spin waves.

*How may these results be modelled?*

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By driving the resonator at the resonance frequency at higher spin wave powers, we observe a sigmoidal-like activation